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Thermodynamically consistent multi-phase DPD

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Mesoscale modelling

Challenge: representing materials at larger-than-atomic scales (10nm and larger, nanoseconds and longer)

- Often requires <u>coarse-graining</u>
 - Collect together atoms/molecules/molecular segments into larger particles ('beads')
 - Work out interactions (potentials, forces) between beads to collectively match up with those between original atoms
- Interactions between beads tend to be soft and (mostly) repulsive
 - Hard-core atomistic interaction types (e.g. Lennard-Jones) not the most realistic at these length and time scales







Dissipative particle dynamics (DPD)

Enhanced molecular dynamics for mesoscale

- Includes additional pairwise forces as momentumconserving thermostat^[1,2]
 - Controls temperature without disrupting hydrodynamics
- Often used in combination with soft interactions between beads to use larger time steps
 - Rapid equilibration compared with atomistic MD
- Has been applied to many soft matter systems involving formation of large-scale structures (e.g. vesicles^[3]) and/or flows^[4]





- 1. Hoogerbrugge and Koelman, *EPL* **19** (3), 155–160 (1992)
- 2. Groot and Warren, *J Chem Phys* **107** (11), 4423–4435 (1997)
- 3. Yamamoto *et al.*, *J Chem Phys* **116** (13), 5842–5849 (2002)
- 4. Posel et al., Soft Matter 13 (8), 1634–1645 (2017)

DPD bead interactions

Default: soft repulsive (Groot-Warren, 'standard DPD') pairwise potential^[1]

$$U(r_{ij}) = \begin{cases} \frac{1}{2}Ar_c \left(1 - \frac{r_{ij}}{r_c}\right)^2, & r_{ij} < r_c \\ 0, & r_{ij} \ge r_c \end{cases}$$

- Finite at zero separation (bounded), zero beyond cutoff (no long-range corrections)
- Similar to potential obtained by coarse-graining polymer melts^[2]
- Leads to quadratic equation of state with bead density ρ :

 $p \approx \rho k_B T + 0.101 A r_c^4 \rho^2$

- Can vary A between different species to match e.g. required energy of mixing
- Single species cannot separate into rarefied/dense (e.g. vapour/liquid) phases



1. Groot and Warren, J Chem Phys 107 (11), 4423–4435 (1997)

2. Forrest and Suter, *J Chem Phys* **102** (18), 7256–7266 (1995)

DPD with improved thermodynamics?

Need to add attraction to interactions to get phase separation

- Could do this by using many-body DPD^[1,2] to impose specific equation of state
 - Additional pairwise calculations of localised bead densities
 - Harder to vary interactions between different species^[3,4]
 - Not easy to ensure thermodynamic stability
- Alternatively, could borrow approach from Mie potential^[5]:

$$U(r_{ij}) = \left(\frac{n}{n-m}\right) \left(\frac{n}{m}\right)^{\frac{m}{n-m}} \epsilon \left[\left(\frac{\sigma}{r_{ij}}\right)^n - \left(\frac{\sigma}{r_{ij}}\right)^m\right]$$

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- 1. Pagonabarraga and Frenkel, *J Chem Phys* **115** (11), 5015–5026 (2001)
- 2. Warren, *Phys Rev E* **68** (6), 066702 (2003)

Control repulsion with n, attraction with m: standard Lennard-Jones interaction when n = 12 and m = 6

- 3. Warren, *Phys Rev E* **87** (4), 045303 (2013)
- 4. Vanya and Elliott, *Phys Rev E* **102** (1), 013312 (2020)
- 5. Mie, Ann Phys **316**, 657–697 (1903)





*n*DPD: extended Groot-Warren interactions^[1]

• Force and potential for bead pairs within cutoff distance *r_c*:

$$\mathbf{F}^{C}(\mathbf{r}_{ij}) = A \left[b \left(1 - \frac{r_{ij}}{r_{c}} \right)^{n} - \left(1 - \frac{r_{ij}}{r_{c}} \right) \right] \frac{\mathbf{r}_{ij}}{r_{ij}}$$
$$U(r_{ij}) = \frac{Abr_{c}}{n+1} \left(1 - \frac{r_{ij}}{r_{c}} \right)^{n+1} - \frac{Ar_{c}}{2} \left(1 - \frac{r_{ij}}{r_{c}} \right)^{2}$$

- Includes form of original interaction as attractive part (similar to many-body DPD model^[2])
- Additional control on repulsion with b and n
- Reduces to 'standard DPD' when b = 2, n = 1





Length scale defined as distance where U = 0 $(b > \frac{n+1}{2}, n \neq 1)$: $\frac{\sigma}{r_c} = 1 - \left(\frac{n+1}{2b}\right)^{\frac{1}{n-1}}$

- 1. Sokhan et al., Soft Matter 19 (30), 5824–5834 (2023)
- 2. Warren, *Phys Rev E* **68** (6), 066702 (2003)

Simulations to explore phase behaviour

Modified DL_MESO_DPD^[1] and DL_POLY^[2] to include new nDPD potential/force

- Used elongated boxes to favour separation into vapour and liquid, to determine phase densities, pressure and surface tension at given temperatures
- Found solid phase structures by annealing (heating and cooling around melting point)
- Simulations confirmed theoretically-determined condition for thermodynamic stability^[3]:

$$b > \frac{(n+1)(n+2)(n+3)(n+4)}{120}$$





n	A	b	σ/r_c	T_c	<i>p</i> _c	ρ_c
2	25.0	3.02	0.5033	1.025	0.2951	0.519
3	15.0	7.2	0.4730	1.284	0.3979	0.504
4	10.0	15.0	0.4497	1.286	0.4095	0.484

1. Seaton, *Mol Sim* **47** (2–3), 228–247 (2021)

- 2. Guest et al., Mol Sim 47 (2–3), 194–227 (2021)
- 3. Louis et al., Phys Rev E 62 (6), 7961–7972 (2000)

Vapour/liquid phase diagram



Plots of temperature against vapour and liquid densities (rescaled by critical values)

- Consistent vapour branch for all n
- Concavity of liquid branch dependent on *n* and *b*
- Liquid branches confirmed for T < 0.7T_c using HNC with nDPD potential^[1] (circles)
- RDFs indicate whether a phase is vapour or liquid
 - 1. https://github.com/patrickbwarren/SunlightHNC

Solid/liquid transitions

Interesting behaviours at low temperatures

- Low melting point $(T_0 \approx 0.07T_c)$
- FCC solid favoured over BCC
- Temperature of maximum (liquid) density (TMD) slightly above melting point
- Solid contracts when heated: negative thermal expansion
- Increasing pressure reduces melting
 point: pressure-induced melting





Simplicity of nDPD interactions belies their ability to model richly complex thermodynamics: many of above effects likely to originate from medium-range (mesoscopic) structure rather than e.g. hydrogen bonding

Parameterising *n*DPD for simple liquid



Variation of saturated vapour pressure (inset: surface tension) with temperature scaled to melting point T_0 , $T' = \frac{T-T_0}{T_c-T_0}$ (compared with reference equation of state^[1])



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Assuming single beads for e.g. water

- Start by setting n (e.g. n = 4)
- Energy scale set by *A*, based on temperature we wish to use compared to *T_c*
- Length scale set by r_c , based on coarsegraining level (molecules per bead)
 - 5 H₂O molecules per bead, $r_c \approx 13.88$ Å
- Value of *b* sets thermodynamic behaviour, can be based on variation of surface tension with temperature

Fitting *n*DPD parameters from coarse-graining

CG/DPD model for *cis*-1,4-polybutadiene^[1]

- Applied trajectory matching technique from MD simulations of cPB with united atom force field
- Devised van der Waals (non-bonded) pairwise forces and potentials for various CG levels
 - λ monomers of 4 C atoms per bead
 - Bond and angle potentials also applied between beads
- Originally fitted CG forces to piecewise polynomial functions of *r* for DPD calculations
- Force/potential looks rather like *n*DPD ...





Fitting *n*DPD for CG polymer

Two CG models available for 5-monomer cPB: fitted *n*DPD parameters using least squares regression

- One based on best likelihood to fit MD trajectories (via Bayesian matching)
 - $b \approx 22.372$ smaller than minimum (23.022) for thermodynamic stability based on n = 4.836
- Other based on best match to actual cPB densities
 - $b \approx 15.949$ thermodynamically stable (n = 4.186)
 - Matched pressure and RDFs using fitted *n*DPD interactions in DL_MESO_DPD and DL_POLY NVT calculations (cf. LAMMPS from original study)







Conclusions

New interaction model (*n*DPD) for CG/mesoscale modelling

- Added attraction and more control on repulsion to Groot-Warren ('standard DPD') interactions
- Enables vapour/liquid phase coexistence below critical point
- Interesting solid/liquid behaviour around melting point despite simplicity of interaction
- Straightforward to parameterise for simple fluids
- Matches interactions obtained by coarsegraining more complex materials (e.g. polymers)





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